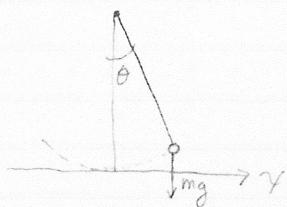


8.)



$$\sin \theta = \frac{x}{l}$$

the potential energy of the mass as a function of  $\theta$ :

$$U(\theta) = mgl(1 - \cos \theta) \quad \text{potential on circular path with } U(0) = 0 \text{ by def}$$

$$U(\theta) = 2mgl(\sin^2 \theta / 2) \quad \text{by trig ident.}$$

taking the derivative with respect to  $\theta$ :

$$dU = 2mg l \sin(\theta/2) \cos(\theta/2) (\frac{1}{2}) d\theta$$

$$dU = mg l \sin \theta d\theta \quad \text{by trig. ident.}$$

dividing by  $dx$

$$\frac{dU}{dx} = mg l \sin \theta \frac{d\theta}{dx}$$

$$\text{substituting } \sin \theta = \frac{x}{l}$$

$$\frac{dU}{dx} = mg l \left(\frac{x}{l}\right) \frac{d}{dx} \left(\sin^{-1} \frac{x}{l}\right)$$

$$\frac{dU}{dx} = mg l \left(\frac{x}{l}\right) \left(\frac{1}{l\sqrt{1-(x/l)^2}}\right)$$

$$\frac{dU}{dx} = mg \left(\frac{x}{l}\right) \left(1 - \left(\frac{x}{l}\right)^2\right)^{-1/2}$$

$$\frac{dU}{dx} = mg \left(\frac{x}{l}\right) \left(1 + \left(\frac{-1}{2}\right)\left(\frac{x^2}{l^2}\right) + \left(\frac{3}{4}\right)\left(\frac{1}{2!}\right)\left(\frac{x^4}{l^4}\right) + \dots\right) \quad \text{by binomial expansion}$$

$$F(x) = -\frac{dU}{dx} \sim -mg \left(\frac{x}{l}\right) \left(1 - \frac{1}{2} \frac{x^2}{l^2}\right)$$

$$F(x) = -mg\left(\frac{x}{l} - \frac{1}{2}\frac{x^3}{l^3}\right)$$

← Force due to the gravitational field  
Now about the centripetal force

set  $F(x) = ma = m\ddot{x}$

$$m\ddot{x} = -mg\left(\frac{x}{l} - \frac{1}{2}\frac{x^3}{l^3}\right)$$

$$\ddot{x} + \left(\frac{g}{l}\right)x - \left(\frac{g}{2l^3}\right)x^3 = 0 \quad \checkmark$$

let  $\omega_0^2 = \frac{g}{l}$  and  $\epsilon = \frac{g}{2l^3}$

$\ddot{x} + \omega_0^2 x - \epsilon x^3 = 0$

the following solution is taken from the book:

To solve this non-linear differential equation,  $\ddot{x} + \omega_0^2 x = \epsilon x^3$ ; assume the solution has the form:

$$x = x_0 + \epsilon x_1$$

Substitute into equation:

$$\ddot{x}_0 + \omega_0^2 x_0 - \epsilon x_0^3 + \epsilon (\ddot{x}_1 + \omega_0^2 x_1 - \epsilon x_1^3) = 0$$

$$\ddot{x}_0 + \omega_0^2 x_0 + \epsilon (\ddot{x}_1 + \omega_0^2 x_1 - x_0^3) + \epsilon^2 (-x_1^3) = 0$$

equating like powers of  $\epsilon$  to zero:

$$\ddot{x}_0 + \omega_0^2 x_0 = 0$$

$$(D^2 + \omega_0^2)x_0 = 0 \Rightarrow \lambda = \pm i\omega_0$$

$$x_0 = A e^{i\omega_0 t} + B e^{-i\omega_0 t}$$

Keeping real-valued functions

$$x_0 = A \cos \omega_0 t + B \sin \omega_0 t$$

$$A = x_0(0) \text{ and } B = \dot{x}_0(0) = 0 \leftarrow \text{choose these initial cond}$$

also have ..

$$\epsilon (\ddot{x}_1 + \omega_0^2 x_1 - x_0^3) = 0$$

$$\ddot{x}_1 + \omega_0^2 x_1 = x_0^3$$

$$\ddot{x}_1 + \omega_0^2 x_1 = A^3 \cos^3 \omega_0 t$$

$$\ddot{x}_1 + \omega_0^2 x_1 = \frac{1}{4} A^3 (3 \cos \omega_0 t + \cos 3\omega_0 t) \quad \text{by trig. ident.}$$

the solution is :

$$x_1(t) = \frac{A^3}{32 \omega_0^2} (12 \omega_0 t \sin \omega_0 t - \cos 3\omega_0 t)$$

now to get rid of those terrible secular terms!

start with the equation of motion:

$$\ddot{x} + \omega_0^2 x - \epsilon x^3 = 0$$

expanding the solution

$$x(t) = x_0 + \epsilon x_1 + \epsilon^2 x_2 + \epsilon^3 x_3 + \dots = \sum_{n=0}^{\infty} (\epsilon)^n x_n$$

expand also the square of the frequency:

$$\omega_0^2 = \omega^2 + \alpha_1 \epsilon + \alpha_2 \epsilon^2 + \dots = \omega^2 + \sum_{n=1}^{\infty} \alpha_n (\epsilon)^n$$

now determine  $\alpha_1, \alpha_2, \dots, \alpha_n$  so that the secular terms are eliminated

substitute into the differential equation of motion:

$$\sum_{n=0}^{\infty} (\epsilon)^n \ddot{x}_n + \left( \omega^2 + \sum_{n=1}^{\infty} \alpha_n (\epsilon)^n \right) \left( \sum_{n=0}^{\infty} (\epsilon)^n x_n \right) - \left( \sum_{n=0}^{\infty} (\epsilon)^{n+1} x_n \right) = 0$$

rearranging:

$$\sum_{n=0}^{\infty} (\epsilon)^n [\ddot{x}_n + \omega^2 x_n] - \sum_{n=0}^{\infty} (\epsilon)^{n+1} x_n + \sum_{n=1}^{\infty} \alpha_n(\epsilon)^n \sum_{r=0}^{\infty} (\epsilon)^r x_n = 0$$

equate like powers of  $\epsilon$  to zero:

$$\ddot{x}_0 + \omega^2 x_0 = 0$$

$$\ddot{x}_1 + \omega^2 x_1 = x_0^3 - \alpha_1 x_0$$

$$\ddot{x}_2 + \omega^2 x_2 = 3x_0^2 x_1 - \alpha_1 x_1 - \alpha_2 x_0$$

choose initial conditions:

$$x(0) = A \quad \dot{x}(0) = 0$$

$$x_0(0) = A \quad \dot{x}_0(0) = 0$$

$$x_i(0) = 0 \quad \dot{x}_i(0) = 0 \quad \text{for } i = 1, 2, 3, \dots, n$$

since  $x_0 = A \cos \omega t$

Substitute into  $\ddot{x}_1 + \omega^2 x_1 = x_0^3 - \alpha_1 x_0$

$$\ddot{x}_1 + \omega^2 x_1 = A^3 \cos^3 \omega t - \alpha_1 A \cos \omega t$$

$$\ddot{x}_1 + \omega^2 x_1 = \left(\frac{3}{4} A^3 - \alpha_1 A\right) \cos \omega t + \frac{1}{4} A^3 \cos 3\omega t$$

since the coefficient of  $\cos \omega t$  must be zero  
since it is a soln. of the homogeneous eq.

$$\alpha_1 = \frac{3}{4} A^2$$

$$\ddot{x}_1 + \omega^2 x_1 = \frac{1}{4} A^3 \cos 3\omega t$$

$$\ddot{x}_1 + \omega^2 x_1 = \frac{1}{4} A^3 \cos 3\omega t$$

$$x_1(t) = \frac{A^3}{32\omega^2} (\cos \omega t - \cos 3\omega t) \quad \text{satisfies initial conditions}$$

and

$$\omega_0 = (\omega^2 + \alpha^2)^{1/2} = (\omega^2 + \frac{3}{4} A^2 \epsilon)^{1/2}$$

thus correct to first order:

$$\boxed{\omega \sim \omega_0 \left( 1 - \frac{3A^2}{8\omega_0^2} \epsilon \right)}$$

relates frequency and amplitude ✓

$$\boxed{x(t) = \left( 1 + \frac{\epsilon A^2}{32\omega^2} \right) A \cos \omega t + \left( \frac{-A^3}{32\omega^2} \right) A \cos 3\omega t}$$

first-order solution to Duffing equation

12.) particle of mass  $m$

$$F(x) = m\ddot{x} \begin{cases} +F_0 & x < 0 \\ -F_0 & x > 0 \end{cases} \quad \text{thus} \quad \ddot{x} = \begin{cases} +F_0/m & x < 0 \\ -F_0/m & x > 0 \end{cases}$$

first consider  $x < 0$

$$\ddot{x} = (F_0/m)$$

$$\ddot{x} = \frac{dv}{dt} = \frac{dv}{dx} \frac{dx}{dt} = v \frac{dv}{dx}$$

$$v dv = \left(\frac{F_0}{m}\right) dx$$

$$\frac{v^2}{2} = \frac{F_0}{m} x + C \quad \dot{x}(x = -A) = 0$$

$$0 = \frac{v^2}{2} = \frac{F_0}{m}(-A) = C \quad \therefore C = \frac{F_0 A}{m}$$

$$\frac{v^2}{2} = \frac{F_0}{m}(x) + \frac{F_0 A}{m}$$

$$v^2 = \frac{2F_0}{m}(x + A) \quad \text{for } x < 0$$

now consider  $x > 0$

$$\ddot{x} = (-F_0/m)$$

$$\frac{v^2}{2} = -\frac{F_0}{m}x + C' \quad \dot{x}(x = +A) = 0$$

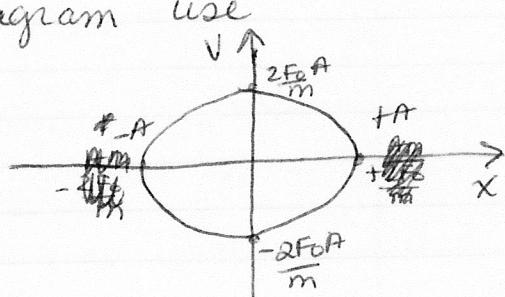
$$0 = -\frac{F_0}{m}A + C' \quad \therefore C' = \frac{F_0}{m}A$$

$$v^2 = \frac{2F_0}{m}(-x + A)$$

thus to graph the phase diagram use

$$\begin{cases} v^2 = \frac{2F_0}{m}(x + A) & \text{for } x < 0 \\ v^2 = \frac{2F_0}{m}(-x + A) & \text{for } x > 0 \end{cases}$$

assumes that the energy of the particle  $E > u(x) \Rightarrow E > \pm F_0 x$



to calculate the period of motion in terms  
of  $m$ ,  $F_0$ , and the amplitude  $A$   
(disregard damping)

calculate time for particle to go from 0 to  $tA$   
and multiply by 4:

$$\ddot{x} = -\frac{F_0}{m}x \quad x > 0$$

$$\dot{x} = -\frac{F_0}{m}xt + C_1 \quad \dot{x}(0) = 0 \quad \therefore C_1 = 0$$

$$x = -\left(\frac{F_0}{2m}\right)t^2 + C_2 \quad x(0) = A \quad \therefore C_2 = A$$

equation of motion in time is

$$x = A - \frac{F_0}{2m}t^2$$

solving for  $t$

$$t = \left[ \frac{2m}{F_0} (A - x) \right]^{\frac{1}{2}}$$

time to go from 0 to  $A$ :

$$\Delta t = \left| \left[ \frac{2m}{F_0} (A - 0) \right]^{\frac{1}{2}} - \left[ \frac{2m}{F_0} (A - A) \right]^{\frac{1}{2}} \right|$$

$$\Delta t = \left( \frac{2mA}{F_0} \right)^{\frac{1}{2}}$$

thus the total period to go all around  
is  $\Delta t(0 \rightarrow A) + \Delta t(A \rightarrow 0) + \Delta t(0 \rightarrow -A) + \Delta t(-A \rightarrow 0)$

$$T = 4\Delta t = 4 \left( \frac{2mA}{F_0} \right)^{\frac{1}{2}} = \left( \frac{32mA}{F_0} \right)^{\frac{1}{2}}$$